

The Effects of Viscosity and Heat Conduction On Sound Waves

Introduction

- The effects of viscosity and heat conduction are not typically considered.
- Consider the damping of the sound waves by analyzing **dispersion relation**.
- Derive the dispersion relation and the damping rate for a non-ideal compressible fluid.
- Heating of the fluid due to viscosity.

Mathematical method

$$\omega = \pm c_s k$$

Linear theory of sound waves:

$$p = p_0 + p_1(\vec{r}, t)$$

Static term and dynamic term where the dynamic term can be treated as first order small term

Complex exponential method:

$$p_1(\vec{r}, t) = \text{Re}\{ \hat{p}_1 \exp(i(\vec{k} \cdot \vec{r} - \omega t)) \}$$

Matrix formation:

$$\begin{bmatrix} \rho_0 & -\gamma p_0 & 0 & 0 & 0 \\ 0 & -\omega & \rho_0 k_x & \rho_0 k_y & \rho_0 k_z \\ k_x & 0 & -\rho_0 \omega & 0 & 0 \\ k_y & 0 & 0 & -\rho_0 \omega & 0 \\ k_z & 0 & 0 & 0 & -\rho_0 \omega \end{bmatrix} \begin{bmatrix} \hat{p}_1 \\ \hat{\rho}_1 \\ u_x \\ u_y \\ u_z \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$

Fourier components – focus on a given wavenumber & frequency (linear theory)

Governing equation

Non-ideal fluid case:

- Navier-Stokes equation:

$$\rho \left(\frac{\partial \vec{u}}{\partial t} + \vec{u} \cdot \nabla \vec{u} \right) = -\nabla p + \mu \nabla^2 \vec{u} + \left(\zeta + \frac{1}{3} \mu \right) \nabla (\nabla \cdot \vec{u})$$

- Continuity equation:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{u}) = 0$$

- Energy equation:

$$\rho C_v \left(\frac{\partial T}{\partial t} + \vec{u} \cdot \nabla T + (\gamma - 1) T \nabla \cdot \vec{u} \right) = \mu \left(\frac{\partial u_i}{\partial x_j} \frac{\partial u_j}{\partial x_i} + \frac{\partial u_i}{\partial x_j} \frac{\partial u_j}{\partial x_i} \right) + \left(\zeta + \frac{1}{3} \mu \right) (\nabla \cdot \vec{u})^2 + K \nabla^2 T$$

- Ideal gas law:

$$p = \frac{k_B}{m} \rho T$$

Ideal fluid case:

- Euler's equation

$$\rho \left(\frac{\partial \vec{u}}{\partial t} + \vec{u} \cdot \nabla \vec{u} \right) = -\nabla p$$

- Continuity equation

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{u}) = 0$$

- Adiabatic energy equation

$$\left(\frac{\partial}{\partial t} + \vec{u} \cdot \nabla \right) \left(\frac{p}{\rho^\gamma} \right) = 0$$

Wave equation:

$$\frac{\partial^2 p_1}{\partial t^2} = c_s^2 \nabla^2 p_1, \text{ where } c_s = \sqrt{\frac{\gamma P_0}{\rho_0}}$$

$$T(x, t) = T_h(x, t) + T_p(x, t)$$

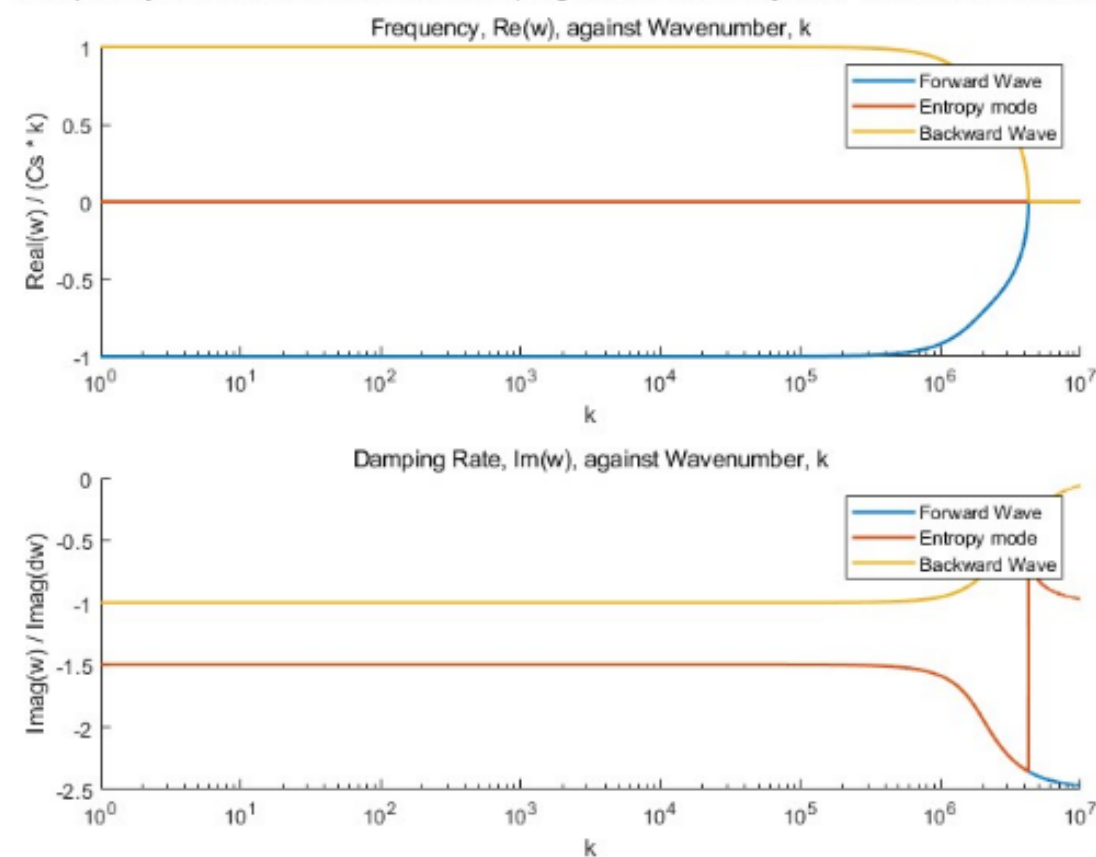
$$T_h(x, t) = \int_{-\infty}^{\infty} \left(\frac{D}{\pi t} \right)^{\frac{1}{2}} \cdot e^{-\frac{(x-y)^2 D}{t}} \cdot A \cos(ky) dy$$

The temperature function (solved by using Green's function)

$$T_p(x, t) = \int_0^t \int_{-\infty}^{\infty} \left(\frac{D}{\pi(t-\tau)} \right)^{\frac{1}{2}} \cdot e^{-\frac{(x-y)^2 D}{t-\tau}} \cdot \{ B \sin(ky - \omega t) + C [\sin(ky - \omega t)]^2 \} dy d\tau$$

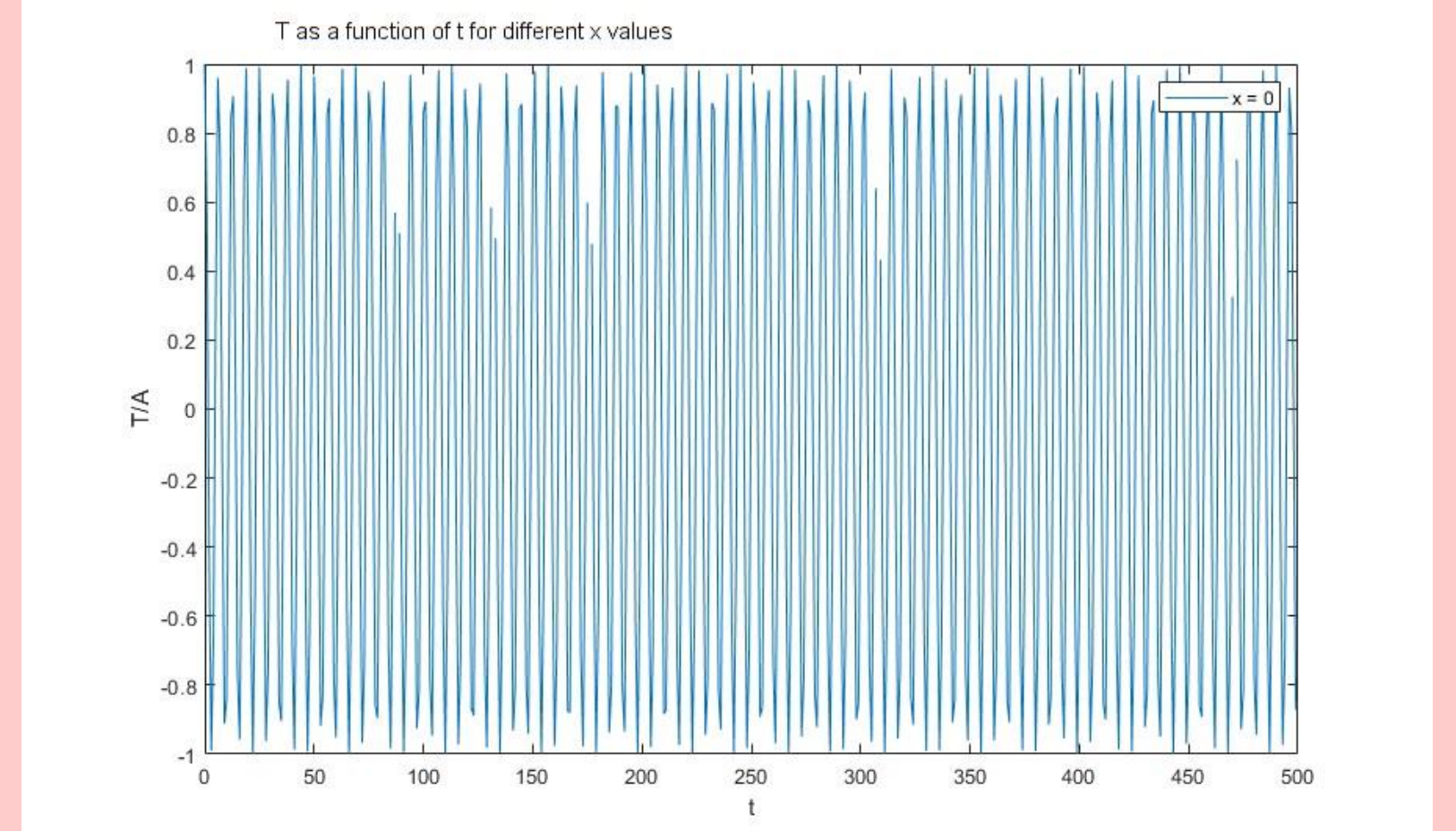
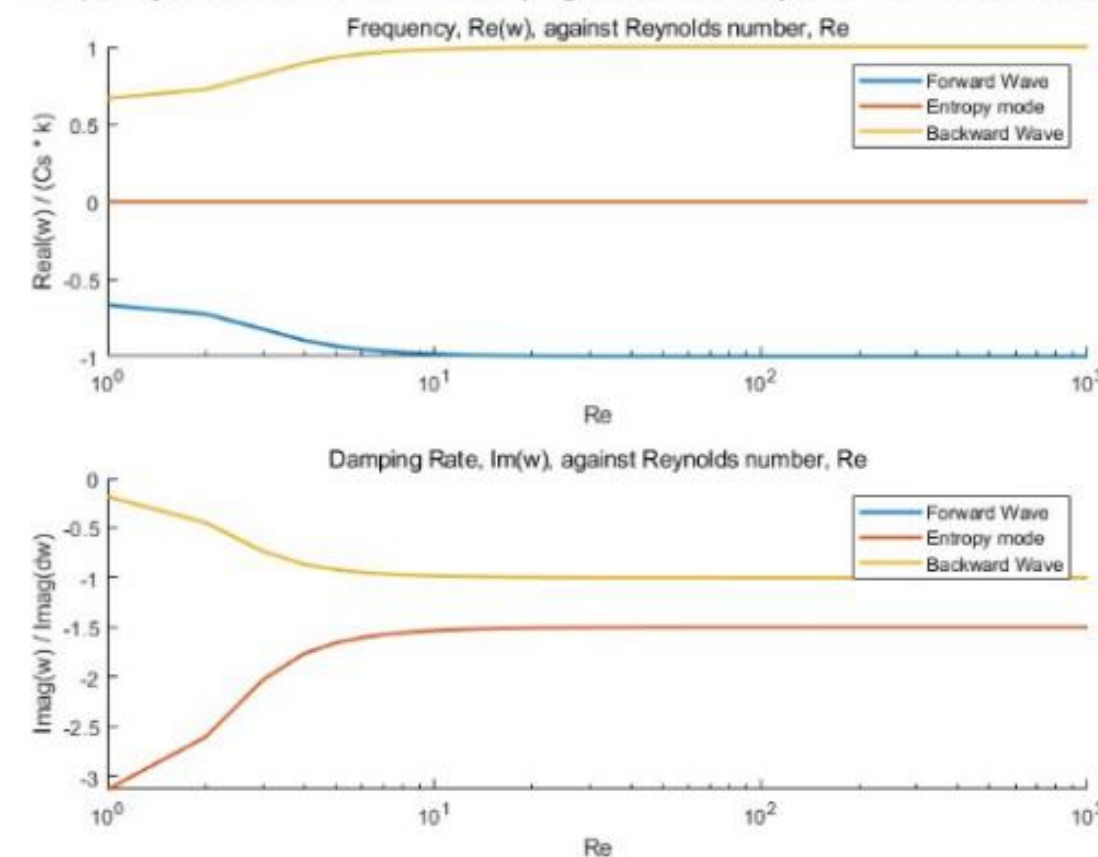
Numerical results

Frequency of Sound Waves with Damping due to Viscosity and Thermal Conduction

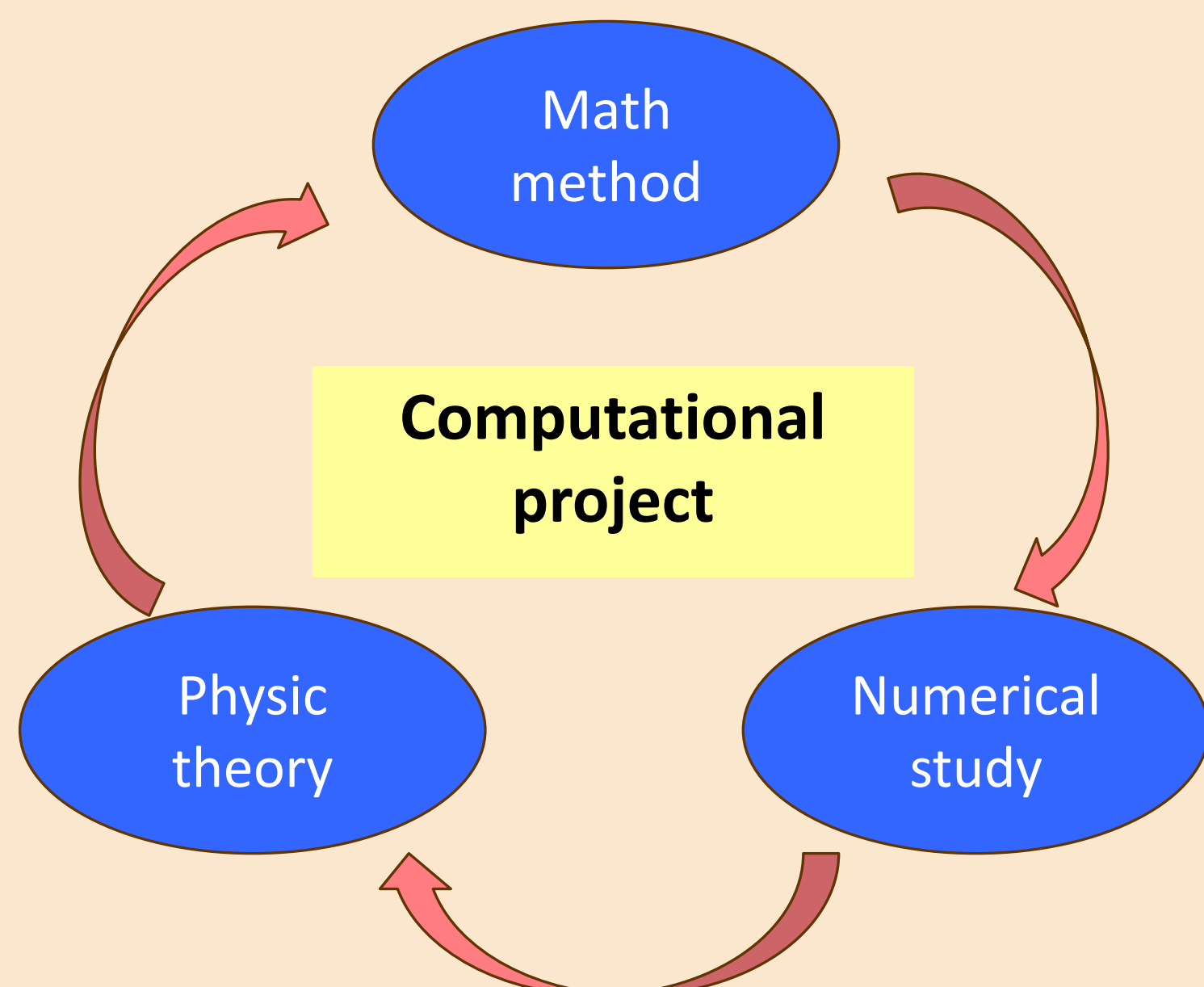


Frequency change for Helium with damping rate with respect of k (left) and Re (right)

Frequency of Sound Waves with Damping due to Viscosity and Thermal Conduction



Temperature change for air with wave compression and expansion



Summary

What we have done:

- Studied the effect of viscosity and heat conduction – dispersion relation
- Heating of gas due viscosity – Green's function solution
- Numerical results

Future investigation:

- Investigate other scenarios where heating could occur e.g. industrial or astrophysical conditions
- Full coupled problem of viscous heating – changes to temperature influence the velocity field
- Reduce numerical error in code